

## GROUP VELOCITY

Group of waves with possibly different phase velocities:

1. **Pure sinusoidal wave:**  $Y(x, t) = A \cos(kx - \omega t + \theta_0) = A \cos(\theta)$

where  $\theta = \theta(x, t) = \text{phase angle} = kx - \omega t + \theta_0$

where  $k = 2\pi \text{ radians}/\lambda$  ( $kx$  is then an angle)

where  $\omega = 2\pi \text{ radians}/T$  ( $\omega t$  is then an angle)

and  $\theta_0$  is the initial phase angle (at  $x=0$  when  $t=0$ )

[ $\theta_0 = 0^\circ$  for cosine wave;  $\theta_0 = -90^\circ$  for sine wave.]

For crest of wave (phase angle of crest is constant at  $90^\circ = \frac{1}{2}\pi$  rad):

$$\theta_{\text{crest}} = \frac{1}{2}\pi = kx_{\text{crest}} - \omega t_{\text{crest}} + \theta_0, \quad \text{or}$$

$$x_{\text{crest}} = (\frac{1}{2}\pi + \omega t_{\text{crest}} - \theta_0) / k .$$

For speed of crest of wave:  $v_{\text{phase}} = v_{\text{crest}} = dx_{\text{crest}}/dt_{\text{crest}}$ , so

$$v_{\text{phase}} = d([\frac{1}{2}\pi + \omega t - \theta_0] / k) / dt = \omega / k \quad (\text{here } t = t_{\text{crest}})$$

for short); recall that  $k=2\pi/\lambda$  and  $\omega=2\pi/T$  so that:

$$v_{\text{phase}} = (2\pi/T) / (2\pi/\lambda) = \lambda/T, \quad \text{and since } f = (1/T)$$

$$\mathbf{v_{\text{phase}} = \omega/k = \lambda f .}$$

2. **Group of waves:** A group of sine waves will add together to form some pattern that also repeats (this is the Fourier Series in reverse).

$$Y_{\text{group}}(x, t) = A(x) \cos(Kx - \omega_g t) \quad \text{where}$$

$A(x)$  is the shape of the group,

$K = 2\pi/\lambda_g$  where  $\lambda_g$  is the distance over which the pattern for the group repeats, and

$\omega_g = 2\pi/T_g$  where  $T_g$  is the time over which the pattern for the group repeats.

At  $t=0$  sec,  $Y_{\text{group}}(x, 0) = A(x) \cos(Kx)$  where  $A(x) = \sum_n a_n \cos(k_n x)$

(here  $A(x)$  is expressed as a Fourier Series), so

$$Y_{\text{group}}(x, 0) = \sum_n a_n \cos(k_n x) \cos(Kx) .$$

We can now use two trig identities

$$[\cos(\theta \pm \phi) = \cos\theta \cos\phi \mp \sin\theta \sin\phi]$$

to get  $\cos\theta \cos\phi = \frac{1}{2}[\cos(\theta + \phi) + \cos(\theta - \phi)]$ , ,

and with  $\theta = k_n x$  and  $\phi = Kx$ , we get

$$Y_{\text{group}}(x, 0) = \sum_n \frac{1}{2} a_n \{ \cos[(k_n + K)x] + \cos[(k_n - K)x] \} ,$$

and since  $\cos(-\theta) = \cos(+\theta)$ , we can write:

$$Y_{\text{group}}(x, 0) = \sum_n \frac{1}{2} a_n \{ \cos[(K + k_n)x] + \cos[(K - k_n)x] \} .$$

**Now put in the time dependence** such that wherever we had a  $Kx$ , we put in an additional  $-\omega t$ :  $(Kx) \rightarrow (Kx - \omega t)$ :

$$Y_{\text{group}}(x, t) = \sum_n \frac{1}{2} a_n \{ \cos[(K+k_n)x - \omega_+ t] + \cos[(K-k_n)x - \omega_- t] \}$$

where we use  $\omega_{\pm}$  to indicate that  $\omega$  depends on  $k=(K \pm k_n)$  .

[Recall that  $v_{\text{phase}} = \omega/k$ , and  $v_{\text{phase}}$  may not be constant but may depend on (vary with)  $\omega$ .] Since  $\omega$  is a function of  $k$  [ $\omega(k) = v_{\text{phase}}k$ ], we can expand  $\omega(k)$  in a Taylor Series about  $k=K$ : normally the "carrier wave" is a high frequency wave with small wavelength and so large  $K$ , and the "signal" is a set of lower frequency waves with larger wavelength values and so smaller  $k$  values.

$$\omega(K \pm k_n) = \omega(K) \pm (d\omega/dk)_K k_n + \text{higher order terms which we neglect ;}$$

$$\text{now let's let } \mathbf{v_g} \equiv (d\omega/dk)_K \text{ so that } \omega_{\pm} = \omega(K \pm k_n) \approx \omega(K) \pm v_g k_n ,$$

so

$$Y_{\text{group}}(x, t) = \sum_n \frac{1}{2} a_n \{ \cos[(K+k_n)x - (\omega + v_g k_n)t] + \cos[(K-k_n)x - (\omega - v_g k_n)t] \}$$

or re-grouping terms:

$$Y_{\text{group}}(x, t) = \sum_n \frac{1}{2} a_n \{ \cos[(Kx - \omega t) + k_n(x - v_g t)] + \cos[(Kx - \omega t) - k_n(x - v_g t)] \} .$$

We can again use our trig identity:

$$\cos(\theta + \phi) + \cos(\theta - \phi) = 2 \cos\theta \cos\phi ,$$

where  $\theta = (Kx - \omega t)$  and  $\phi = k_n(x - v_g t)$  , to get:

$$Y_{\text{group}} = \sum_n a_n \cos(Kx - \omega t) \cos[k_n(x - v_g t)] ;$$

but here the  $\cos(Kx - \omega t)$  can come out of the summation, so

$$\mathbf{Y_{group}} = \{ \sum_n a_n \cos[k_n(x - v_g t)] \} \cos(Kx - \omega t) = \mathbf{A(x - v_g t) \cos(Kx - \omega t)},$$

where we identify the function  $A(x - v_g t)$  as the original Fourier series with  $x$  replaced by  $(x - v_g t)$ ; that is, the shape moves through space with a speed of  $v_g$ , hence the name **group velocity**.

### 3. Review:

$$\mathbf{v_{phase} = \omega/k = \lambda f}$$

(good for any pure sine wave [or cosine wave] of wavelength  $\lambda$  and frequency  $f$  [or wavevector  $k$  and angular speed  $\omega$ ]) ;

$$\mathbf{v_{group} = d\omega/dk} .$$

### 4. Special case: If $v_{\text{phase}} = \text{constant}$ , then $\omega = v_{\text{phase}}k$ , and so

$$v_{\text{group}} = d\omega/dk = d[v_{\text{phase}}k]/dk = v_{\text{phase}} .$$